Ketten bilden über gemeinsame Kanten aus Sauerstoffatomen auch in der $y$-Richtung unbegrenzte Verbände. Benachbarte zweidimensionale Netzwerke dieser Art werden durch die Anionen und F-H...O-Wasserstoffbrücken miteinander verbunden.

Die Autoren danken Herrn Privatdozent Dr H. Falius, Braunschweig, für die Anregung zu dieser Arbeit und Überlassung von Kristallmaterial, dem Rechenzentrum der Technischen Universität Braunschweig und dem Deutschen Rechenzentrum in Darmstadt für Rechenzeit sowie der Deutschen Forschungsgemeinschaft und dem Fonds der Chemischen Industrie für Leihgaben und Sachmittel.

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Acta Cryst. (1971): B27, 1986

# The Crystal Structure of 4-Cyanopyridine 

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(Received 30 October 1970)


#### Abstract

The crystal structure of 4-cyanopyridine, $\mathrm{C}_{6} \mathrm{~N}_{2} \mathrm{H}_{4}$, has been determined from three-dimensional singlecrystal X-ray data collected by standard film techniques. The crystals are orthorhombic, space group $P c c n$, with $a=3.82 \pm 0 \cdot 01, b=6.12 \pm 0.01, c=23 \cdot 22 \pm 0.02 \AA ; Z=4 ; D_{\text {obs }}=1 \cdot 27, D_{\text {calc }}=1 \cdot 27 \mathrm{~g}^{2} \mathrm{~cm}^{-3}$. The structure was refined by the least-squares method to a conventional $R=0 \cdot 105$ for 550 observed and unobserved data. The molecule retains the twofold axis through the cyano group and pyridine nitrogen atom as a required symmetry element in the unit cell. The bond lengths and angles are quite normal: The mode of packing is unusual, and results in a relatively large 'hole' running through the crystal parallel to $\mathbf{a}$.


## Introduction

4-Cyanopyridine is unique among the simpler 4 -substituted derivatives of pyridine, as at room temperature it is a white crystalline solid with little odour, while for example, the 4-methyl, ethyl, vinyl, and chloro compounds are liquids, foul smelling and easily oxidized (Heilbron, 1965). The mode of packing of the molecules was therefore of interest and the structure determination was undertaken.

## Experimental

Material was purchased from Aldrich Chemical Company and suitable crystals were obtained by sublimation. A well-formed rectangular parallelepiped $0.5 \times$ $0.2 \times 0.2 \mathrm{~mm}$ was sealed in a capillary and used for all subsequent X-ray work. Oscillation, Weissenberg and precession photographs showed the unit cell to be orthorhombic, and the systematic extinctions, 0 kl for $l$ odd, $h 0 l$ for $l$ odd and $h k 0$ for $h+k$ odd, were consistent with space group Pccn. Unit-cell dimensions were $a=3 \cdot 82 \pm 0 \cdot 01, b=6 \cdot 12 \pm 0 \cdot 01, c=23 \cdot 22 \pm 0 \cdot 02 \AA$, measured from calibrated $0 \bar{k} l$ Weissenberg and $h 0 l$ precession photographs. The density was measured by flota-
tion in a chlorobenzene-bromobenzene mixture to be $1.27 \pm 0.01 \mathrm{g.cm}^{-3}$, showing 4 molecules per unit cell.
Because there are 8 equivalent positions in the space group, and only 4 molecules per unit cell, the asymmetric unit consists of only half of the molecule. The twofold axis passing through the cyano group and the nitrogen and $\gamma$ carbon atoms of the pyridine ring must therefore be coincident with a twofold axis parallel to $\mathbf{c}$ in the unit cell. The $x, y$ coordinates of these four atoms are thus required to be $0 \cdot 25,0 \cdot 25$.
Integrated multiple-film data were collected by the equi-inclination Weissenberg method using Ni-filtered $\mathrm{Cu} K \alpha$ radiation for layers 0 kl to 3 kl . Of the 550 recordable data, 450 were observed above background. The intensities were estimated by visual comparison with a time-calibrated intensity strip. Lorentz and polarization corrections were calculated and applied to the data by hand. Precession photographs were taken of layers $h 0 l$ to $h 2 l$ and these data were used to scale the Weissenberg data and to supply intensities for certain reflections that had been obscured by the beam catcher. Unobserved data were given the value local $I_{\min } / 2$.
A three-dimensional temperature-sharpened Patterson map was calculated using the 450 observed data. (The values of $F_{\text {obs }}^{2}$ were multiplied by the value


Fig. 1. Bond lengths and angles. The estimated standard deviations in bond lengths are $0.006 \AA$ for $\mathrm{N}(1)-\mathrm{C}(1)$, $C(1)-C(2)$ and $C(2)-C(3)$, and $0.008 \AA$ for $C(3)-C(4)$ and $\mathrm{C}(4)-\mathrm{N}(2)$, and in bond angles are $0 \cdot 3^{\circ}$ for $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{C}(4)$ and $0.5^{\circ}$ for all others involving carbon and nitrogen atoms. The bond lengths have not been corrected for thermal motion.
$\exp (8.0 / S)$ where $S$ was the appropriate value of $\left.\sin ^{2} \theta / \lambda^{2}\right)$. The pyridine ring-pyridine ring vector of sixfold weight corresponded to the largest peak on the line $u=0 \cdot 5, v=0.5$. From the $w$ coordinate of this peak was deduced the $z$ coordinate of the approximate midpoint of the pyridine ring. The relative coordinates of the atoms within the molecule were easily obtained from the well-defined peaks near the origin, while the orientation of the cyano group was deduced from the relative heights and positions of the peaks in the Harker section at $u=0 \cdot 5$. The approximate $x, y, z$ coordinates for all the atoms were thus directly determined (Sparrow, 1970).

The structure was refined isotropically to $R=0.2$ by a series of observed and difference Fourier summations, the coordinates being adjusted by the method of Booth (1948). The contributions of the two unique hydrogen atoms were clearly visible in both observed and difference Fourier maps. The two carbon atoms $\mathrm{C}(1)$ and $\mathrm{C}(2)$ in the general positions (see Fig. 1) were somewhat elongated parallel to the $x$ direction, showing anisotropic thermal motion. All calculations to this stage were carried out on an 8 k core IBM 1130 com-
puter at the University of Natal, with the local set of programs (Laing \& Sale, 1969).

The structure was then refined anisotropically by the method of least squares using the CRYM system at the California Institute of Technology. A Hughes (1941) weighting scheme was used. Weights were proportional to $1 / F_{\text {obs }}^{2}$ for $F_{\text {obs }}$ greater than $4 F_{\text {min }}$ and to $1 /\left(F_{\text {ots }} \times\right.$ $4 F_{\text {min }}$ ) for $F_{\text {obs }}$ less than $4 F_{\text {min }}$, which was set at $2 \cdot 5$. The function minimized was $\sum w\left(k^{2}\left|F_{\text {obs }}\right|^{2}-\left|F_{\text {calc }}\right|^{2}\right)^{2}$ where $w=$ weight, $k=$ scale factor.

Neutral atom scattering factors for $\mathrm{N}, \mathrm{C}$, and H were taken from International Tables for X-ray Crystallography (1962). An extinction correction of $20 \times 10^{-6}$ was applied (Larson, 1967), but this parameter was not included in the refinement; absorption corrections were not applied. The hydrogen atoms were included in the structure factor calculations (with positional coordinates deduced from the observed and difference Fourier maps), but their parameters were not refined. The final value of the conventional $R$ was $0 \cdot 105$ fur the $550 \mathrm{ob}-$ served and unobserved data. The final positional and thermal parameters are given in Table 1. The observed and calculated structure factors are listed in Table 2.
Intermolecular non-bonded separations are given in Table 3. Fig. 1 shows the numbering system and the bond lengths and angles. Fig. 2 shows the electron density map, calculated in the plane of the molecule. The thermal vibration ellipsoids are illustrated in Fig. 3, and their parameters are given in Table 4. The mode of packing of the molecules and certain intermolecular separations are shown in Fig. 4.

## Discussion

The molecule is planar, with normal bond lengths and angles. The $\mathrm{C}-\mathrm{N}$ bond length of $1.331 \AA$ in the pyridine ring may be compared with the values of $1.335 \AA$ in isoniazide (Kupfer \& Tsoucaris, 1963), $1 \cdot 340 \AA$ in pyridine (Bak, Hansen-Nygaard \& Rastrup-Andersen, 1958), $1.338 \AA$ in sym-triazine (Lancaster \& Stoicheff, 1956) and $1.33 \AA$ in pyridine- $N$-oxide hydrochloride (Tsoucaris, 1961). The $\mathrm{C}-\mathrm{C}$ bond lengths of 1.381 and $1.383 \AA$ in the ring are close to those observed in isonia-

Table 1. Positional parameters, anisotropic temperature factors, and isotropic temperature factors* The general form of the anisotropic temperature factor is $\exp \left[-\left(\beta_{11} h^{2}+\beta_{22} k^{2}+\beta_{33} l^{2}+\beta_{12} h k+\beta_{13} h l+\beta_{23} k l\right)\right]$.


* The estimated standard deviations, obtained from the $C R Y M$ structure-factor least-squares program, are shown in parentheses. A full-matrix least-squares refinement by one of the referees indicates that these values are too small, and that e.s.d's of $t$ twice these values are more realistic. The approximate errors in the hydrogen atom parameters are $x \pm 0.02, y \pm 0.01, z \pm 0.005$, $B \pm 1 \cdot 0$.

Table 2. Observed and calculated structure factors ( $\times 10$ ).
The three columns are $l$ index, $F_{\text {obs }}, F_{\text {calc }}$. Unobserved data are marked with asterisks.


Table 3. Intermolecular non-bonded separations
$\left.\begin{array}{ll}\begin{array}{c}\text { Atom in } \\ \text { asymmetric to } \\ \text { unit }\end{array} \\ \text { atom }\end{array}\right)$

Equivalent positions

| $\prime$ | $x+\frac{1}{2},-y-1,-z+\frac{1}{2}$ |
| :--- | :--- |
| $\prime \prime$ | $x+\frac{1}{2}$, |
| $\prime \prime$ | $y+\frac{1}{2},-z$ |
| $\prime$ | $-x+\frac{3}{2}$, |
|  | $-y+\frac{1}{2}, \quad z$ |

The least-squares plane is defined by atoms $\mathbf{C}(1) \mathrm{C}(2) \mathrm{C}(3)$ $\mathrm{C}(4) \mathrm{N}(1) \mathrm{N}(2)$. Direction cosines of the normal to the plane referred to $a, b$ and $c$ are $0.930,-0.367$ and 0.000 respectively. Origin-to-plane distance $=0.33 \AA$
Deviation from plane $(\AA)$ :

| $C(1)$ | $C(2)$ | $C(3)$ | $C(4)$ | $\mathrm{N}(1)$ | $\mathrm{N}(2)$ | $\mathrm{H}(1)$ | $\mathrm{H}(2)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.004 | -0.004 | 0.000 | 0.000 | 0.000 | 0.000 | 0.00 | 0.01 |

Table 4. The observed thermal ellipsoids

|  | $\sqrt{\overline{u^{2}}}$ | Direction cosines referred to |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | (A) | $a$ | $b$ | $c$ |
| C(1) | $0 \cdot 21$ | -0.092 | -0.369 | -0.925 |
|  | $0 \cdot 25$ | $0 \cdot 198$ | -0.917 | $0 \cdot 346$ |
|  | $0 \cdot 28$ | -0.976 | -0.151 | $0 \cdot 158$ |
| C(2) | 0.21 | 0.096 | -0.014 | 0.995 |
|  | $0 \cdot 23$ | $0 \cdot 200$ | $0 \cdot 980$ | -0.006 |
|  | $0 \cdot 25$ | $0 \cdot 975$ | -0.200 | $-0.097$ |
| C(3) | $0 \cdot 20$ | 0.0 | $0 \cdot 0$ | $-1.0$ |
|  | $0 \cdot 22$ | -0.996 | 0.091 | 0.0 |
|  | $0 \cdot 24$ | 0.091 | $0 \cdot 996$ | $0 \cdot 0$ |
| C(4) | $0 \cdot 20$ | $0 \cdot 0$ | $0 \cdot 0$ | $-1.0$ |
|  | $0 \cdot 24$ | -0.589 | -0.808 | $0 \cdot 0$ |
|  | $0 \cdot 29$ | -0.808 | 0.589 | 0.0 |
| N(1) | $0 \cdot 21$ | $0 \cdot 0$ | 0.0 | 1.0 |
|  | $0 \cdot 25$ | -0.239 | $0 \cdot 971$ | $0 \cdot 0$ |
|  | $0 \cdot 30$ | 0.971 | 0.239 | $0 \cdot 0$ |
| N(2) | $0 \cdot 21$ | $0 \cdot 0$ | 0.0 | $1 \cdot 0$ |
|  | $0 \cdot 26$ | $0 \cdot 470$ | $0 \cdot 883$ | $0 \cdot 0$ |
|  | $0 \cdot 40$ | 0.883 | -0.470 | 0.0 |

zide which range from 1.381 to $1.389 \AA$ (Kupfer \& Tscoucaris, 1963) and to those of $1 \cdot 394 \AA$ found in pyridine (Bak et al., 1958).
The values of $1.439 \AA$ for C (ring)-C(nitrile) and $1 \cdot 137 \AA$ for $\mathrm{C} \equiv \mathrm{N}$ are similar to the expected values of 1.43 and $1 \cdot 15 \AA$ respectively. (Sutton, 1965). The


Fig. 2. Electron density map calculated in the plane of the molecule. Contours are drawn at $0 \cdot 5,1,2,3$ etc. e. $\AA^{-3}$.


Fig. 3. Thermal vibration ellipsoids, viewed in projection down b.
$\mathrm{C}-\mathrm{N}-\mathrm{C}$ and $\mathrm{N}-\mathrm{C}-\mathrm{C}$ angles of 117.3 and $123.9^{\circ}$ are close to the values of $116.9^{\circ}$ and $123 \cdot 8^{\circ}$ respectively observed in pyridine.
There are few short intermolecular non-bonded contacts and it appears that the forces holding the structure together are of three types: (i) between the $\pi$ clouds of the pyridine rings stacked by translation along a, (ii) between pairs of $\mathrm{C} \equiv \mathrm{N}$ groups related by centres of symmetry, and (iii) for a molecule of type 1 (Fig. 4), between $\mathrm{C}(1)-\mathrm{H}(1)$ and the pyridine nitrogen $\mathrm{N}(1)$ of a molecule of type 2 (related by a $2_{1}$ screw axis), and between $\mathrm{C}(2)-\mathrm{H}(2)$ and the cyano nitrogen $\mathrm{N}(2)$ of a molecule of type 3 (related by a centre of inversion).

While the two $\mathrm{C}-\mathrm{H} \cdots \mathrm{N}$ arrangements can hardly be called 'hydrogen bonds' (Donohue, 1968), their presence seems to be more than coincidence because they effectively prevent the usual 'bump-in-hollow' packing parallel to b (Kitaigorodskii, 1961).
The distance between $\mathrm{N}(1)$ in any molecule and $\mathrm{N}(2)$ of the molecule related to it by the $c$ glide, is $6 \cdot 27 \AA$. These molecules do not 'touch'; neither do those related by unit translation parallel to $\mathbf{b}$. This results in a 'hole' of diameter about $1 \AA$ that runs through the crystal parallel to a. Its presence is the cause of the relatively low density of the material.

The thermal motion parameters show that the molecule is vibrating about a point close to $C(3)$, with the cyano group 'wagging' roughly parallel to the $a$ axis. This is physically reasonable, as the $\mathrm{N} \cdots \mathrm{N}$ separation in this direction is $3.82 \AA$, while the contacts in other directions are shorter; $3 \cdot 41,3 \cdot 62$, and $3.71 \AA$ (Table 3, Fig. 4) [The relatively high thermal motion of the structure as a whole was evident from the areas of diffuse scattering surrounding certain reflexions on the Weissenberg photographs. This effect is well known in crystals of planar aromatic compounds (Hoppe, 1957; Amoros, Canut \& de Acha, 1960; Laing, 1964).]

It would appear that this structure is not 'maximum close packed' as a result of the molecule retaining its twofold axis of symmetry in the crystal. While space group Pccn is permissible (but improbable for molecules of symmetry $\overline{1}$ ), space group $P b c n$ would be more suitable for molecules of symmetry 2, while Pnma would be more suitable for molecules of symmetry $m$ (Kitaigorodskii, 1961, p. 100). Were the molecule to retain no symmetry element (other than 1) in the crystal, it is easy to show, with the aid of suitable ('Leybold') models, that a more dense structure could be achieved by simple 'bump-in-hollow' packing by translation. One is therefore led to conclude that this is a unique structure, and that the $\mathrm{C}-\mathrm{H} \cdots \mathrm{N}$ interactions somehow force the molecules to pack in this manner (Fig. 5).


Fig. 4. Mode of packing of the molecules, viewed in projection down a.

We are grateful to the South African Council for Scientific and Industrial Research for a grant to cover running expenses and to the Chemistry Department, California Institute of Technology, for a grant to cover the cost of the final refinement. M. Laing wishes to thank the faculty of the Chemistry Department, San Fernando Valley State College and especially Dr R. E. Marsh, of the California Institute of Technology for their generosity and hospitality during his stay in Los Angeles in 1970. The diagrams were prepared by Miss Lillian Casler.

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Fig. 5. Mode of packing of the molecules related by simple translations along a and $\mathbf{b}$, viewed in projection down $\mathbf{c}$.

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# A Further Investigation of the Crystal Structures of Naphthazarin 

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(Received 26 October 1969 and in revised form 25 July 1970)
The crystal structures of two modifications of naphthazarin (formally 5,8-dihydroxy-1,4-naphthaquinone) have been further refined, using three-dimensional X-ray counter data in one case, and previously reported photographic data in the other. The carbon skeleton of the molecule is shown to possess $D_{2 h}$ symmetry within the standard deviation in carbon-carbon bond length, which was $0.006 \AA$ in the more accurate study. The phenolic hydrogen atom, which is involved in an intramolecular hydrogen bond between the phenolic and quinonoid oxygen atoms, appears to be non-symmetrically placed between them. This structure can be interpreted in terms of resonance between the zwitterion contributors:



## Introduction

Although the compound known as naphthazarin is commonly described as 5,8 -dihydroxy-1,4-naphthaquin-
one, there has been some doubt as to its true structure. There is little difference in nuclear arrangement between the equivalent structures (I) and (III), and it seemed possible that these may have been canonicals of a sym-

